

The Three-Loop Splitting Functions in QCD

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with

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Introduction

Notation, Method (\rightarrow parallel session)

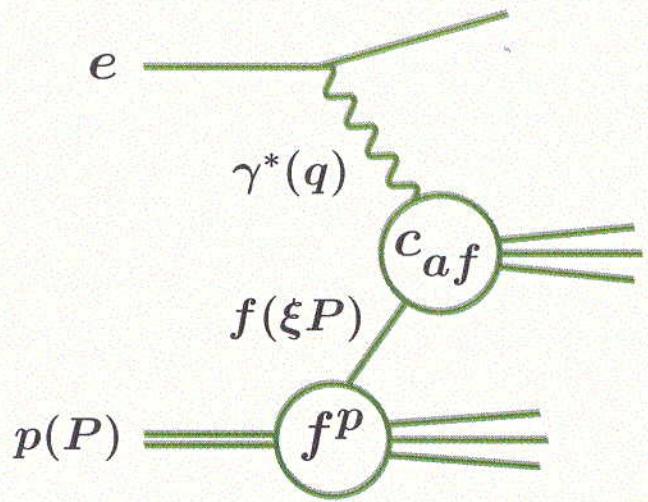
Non-singlet and singlet results

Large x , small x , numerical effects

Summary

hep-ph/0403192 (ns), hep-ph/0404111 (s)

DIS in perturbative QCD



Kinematic variables

$$Q^2 = -q^2$$

$$x = Q^2/(2P \cdot q)$$

Structure functions F_a [up to $\mathcal{O}(1/Q^2)$]

$$\eta_a F_a^p(x, Q^2) = \sum_f \left[c_{a,f}(a_s) \otimes f^p(\mu^2) \right] (\xi)$$

Parton distribution functions f (PDF's)

$$\frac{d}{d \ln \mu^2} f(\xi, \mu^2) = \sum_{f'} \left[P_{ff'}(a_s) \otimes f'(\mu^2) \right] (\xi)$$

Coefficient functions c , splitting functions P

$$P = a_s P^{(0)} + a_s^2 P^{(1)} + a_s^3 P^{(2)} + \dots$$

$$c = \underbrace{a_s^{n_a} \left[c^{(0)} + a_s c^{(1)} + a_s^2 c^{(2)} + \dots \right]}_{\text{NLO:}}$$

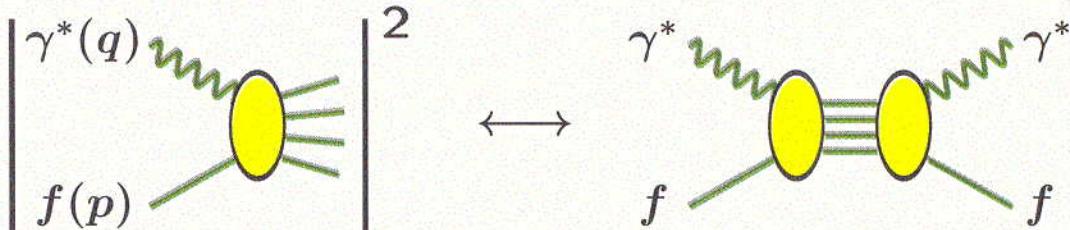
NLO:

Standard approximation for pert. QCD analyses

Next-to-next-to-leading order (NNLO): $c^{(2)}$, $P^{(2)}$

How the calculation is done

Optical theorem: forward Compton amplitude



Coefficient of $(2p \cdot q)^N \leftrightarrow N\text{-th Mellin moment}$

$$A^N = \int_0^1 dx x^{N-1} A(x)$$

P_{gg}, P_{gq} : DIS with scalar ϕ coupling to $G_{\mu\nu}^a G_a^{\mu\nu}$

Gluon polarization sum \leftrightarrow diagrams with ghost h

	tree	1-loop	2-loop	3-loop
$q\gamma$	1	3	25	359
$g\gamma$		2	17	345
$h\gamma$			2	56
qW	1	3	32	589
$q\phi$		1	23	696
$g\phi$	1	8	218	6378
$h\phi$		1	33	1184
sum	3	18	350	9607

Highly optimised symbolic treatment needed

FORM

J. Vermaseren (1989-2004)

Capabilities substantially extended for this project

Evolution of parton distributions

Non-singlet and singlet distributions q^\pm , q^v and q_S , g

$$q_{ik}^\pm = q_i \pm \bar{q}_i - (q_k \pm \bar{q}_k)$$

$$q^v = \sum_{r=1}^{n_f} (q_r - \bar{q}_r)$$

$$f_S = \begin{pmatrix} q_S \\ g \end{pmatrix}, \quad q_S = \sum_{r=1}^{n_f} (q_r + \bar{q}_r)$$

Evolution equations for $\mu_r = \mu_f$

$$\frac{d}{d \ln \mu_f^2} f(x, \mu_f^2) = [P(\alpha_s(\mu_f^2)) \otimes f(\mu_f^2)](x)$$

$2n_f - 1$ scalar (ns) equations + 2×2 singlet case

Mellin convolution

$$[a \otimes b](x) = \int_x^1 d\xi \, a\left(\frac{x}{\xi}\right) b(\xi)$$

Splitting-function combinations

$$P_{\text{ns}}^\pm, \quad P_{\text{ns}}^v = P_{\text{ns}}^- + P_{\text{ns}}^S$$

$$P_S = \begin{pmatrix} P_{qq} & P_{qg} \\ P_{gq} & P_{gg} \end{pmatrix}$$

$$P_{qq} = P_{\text{ns}}^+ + P_{\text{ps}}$$

LO: $P_{\text{ns}}^+ = P_{\text{ns}}^- = P_{\text{ns}}^v = P_{qq}$, NLO: $P_{\text{ns}}^- = P_{\text{ns}}^v$

Basic functions in N - and x -space

N -space: harmonic sums

$$S_{\pm m}(M) = \sum_{i=1}^M \frac{(\pm 1)^i}{i^m}$$

$$S_{\pm m_1, m_2, \dots, m_k}(M) = \sum_{i=1}^M \frac{(\pm 1)^i}{i^{m_1}} S_{m_2, \dots, m_k}(i)$$

Weight = sum of absolute values of indices

Notation: $N_{\pm i} S_{\vec{m}} = S_{\vec{m}}(N \pm i)$

Vermaseren (98)

x -space: harmonic polylogarithms

$$H_0(x) = \ln x, \quad H_{\pm 1}(x) = \mp \ln(1 \mp x)$$

$$H_{m_1, \dots, m_w}(x) = \begin{cases} \frac{1}{w!} \ln^w x, & \vec{m} = \vec{0} \\ \int_0^x dz f_{m_1}(z) H_{m_2, \dots, m_w}(z), & \text{else} \end{cases}$$

with

$$f_0(x) = \frac{1}{x}, \quad f_{\pm 1}(x) = \frac{1}{1 \mp x}$$

Weight = number of indices (0, ± 1)

$$H_{\underbrace{0, \dots, 0}_m, \pm 1, \underbrace{0, \dots, 0}_n, \pm 1, \dots}(x) \equiv H_{\pm(m+1), \pm(n+1), \dots}(x)$$

Remiddi, Vermaseren (99)

$$P_{gg}^{(0)}(x) =$$

$$C_A \left(4[(1-x)^{-1} + x^{-1} - 2 + x - x^2] + \frac{11}{3} \delta(1-x) \right) - \frac{2}{3} n_f \delta(1-x)$$

$$P_{gg}^{(1)}(x) =$$

$$\begin{aligned} & 4C_A n_f \left(1 - x - \frac{10}{9} p_{gg}(x) - \frac{13}{9} \left(\frac{1}{x} - x^2 \right) - \frac{2}{3} (1+x) H_0 - \frac{2}{3} \delta(1-x) \right) + 4C_A^2 \left(27 \right. \\ & + (1+x) \left[\frac{11}{3} H_0 + 8H_{0,0} - \frac{27}{2} \right] + 2p_{gg}(-x) \left[H_{0,0} - 2H_{-1,0} - \zeta_2 \right] - \frac{67}{9} \left(\frac{1}{x} - x^2 \right) \\ & - 12H_0 - \frac{44}{3} x^2 H_0 + 2p_{gg}(x) \left[\frac{67}{18} - \zeta_2 + H_{0,0} + 2H_{1,0} + 2H_2 \right] + \delta(1-x) \left[\frac{8}{3} + 3\zeta_3 \right] \left. \right) \\ & + 4C_F n_f \left(2H_0 + \frac{21}{3x} + \frac{10}{3} x^2 - 12 + (1+x) \left[4 - 5H_0 - 2H_{0,0} \right] - \frac{1}{2} \delta(1-x) \right) \end{aligned}$$

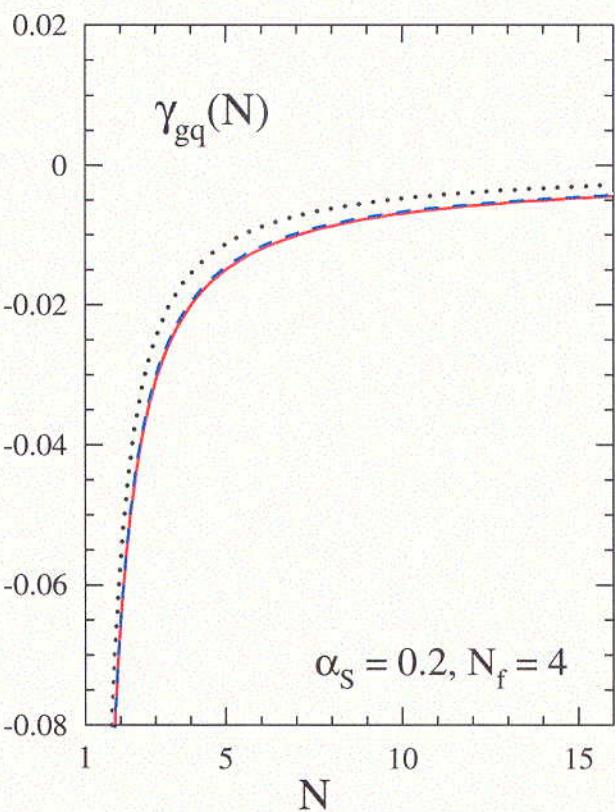
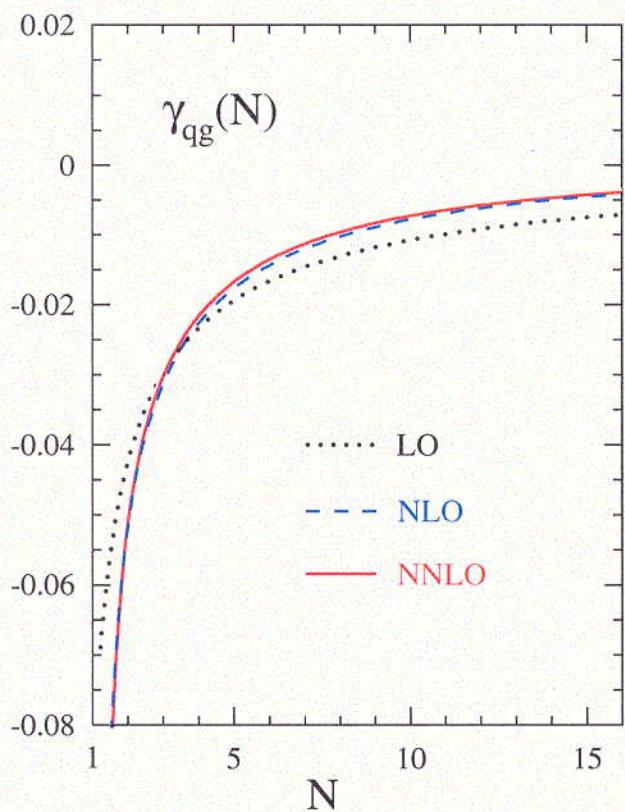
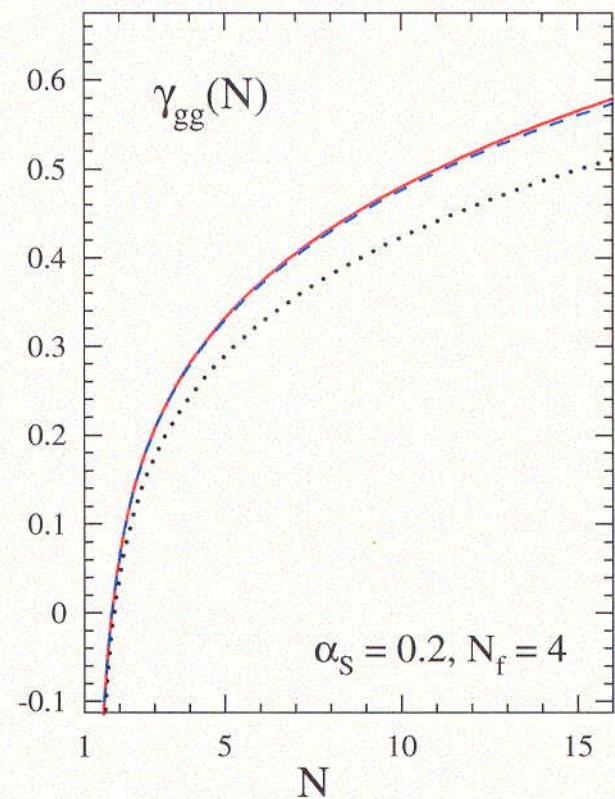
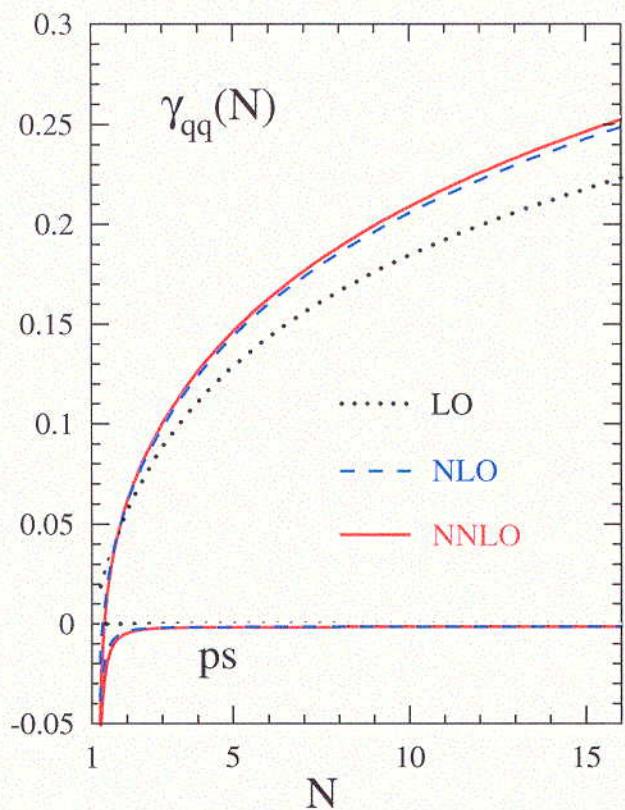
with

$$p_{gg}(x) = (1-x)^{-1} + x^{-1} - 2 + x - x^2$$

$$P_{gg}^{(2)}(x) =$$

$$\begin{aligned}
& 16C_A C_F n_f \left(x^2 \left[\frac{4}{9} H_2 + 3H_{1,0} - \frac{97}{12} H_1 + \frac{8}{3} H_{-2,0} - \frac{2}{3} H_0 \zeta_2 + \frac{103}{27} H_0 - \frac{16}{3} \zeta_2 + 2H_3 \right. \right. \\
& - 6H_{-1,0} + 2H_{2,0} + \frac{127}{18} H_{0,0} - \frac{511}{12} \Big] + p_{gg}(x) \left[2\zeta_3 - \frac{55}{24} \right] + \frac{4}{3} \left(\frac{1}{x} - x^2 \right) \left[\frac{17}{24} H_{1,0} - \frac{43}{18} H_0 \right. \\
& - \frac{521}{144} H_1 - \frac{6923}{432} - \frac{1}{2} H_{2,1} + 2H_1 \zeta_2 + H_0 \zeta_2 - 2H_{1,0,0} + \frac{1}{12} H_{1,1} - H_{1,1,0} - H_{1,1,1} \Big] - \frac{175}{12} H_2 \\
& + 6H_{-1,0} + 8H_0 \zeta_3 - 6H_{-2,0} - \frac{53}{6} H_0 \zeta_2 - \frac{49}{2} H_0 + \frac{185}{4} \zeta_2 + \frac{511}{12} - \frac{1}{2} H_{2,0} - 3H_{1,0} - 4H_{0,0,0,0} \\
& - \frac{67}{12} H_{0,0} + \frac{43}{2} \zeta_3 - H_{2,1} + \frac{97}{12} H_1 - 4\zeta_2^2 - \frac{9}{2} H_3 - 8H_{-3,0} + \frac{33}{2} H_{0,0,0} + \frac{4}{3} \left(\frac{1}{x} + x^2 \right) \left[\frac{1}{2} H_2 - H_{2,0} \right. \\
& + \frac{11}{3} H_{-1,0} + H_{-2,0} + \frac{19}{6} \zeta_2 + 2\zeta_3 - H_{-1} \zeta_2 - 4H_{-1,-1,0} - \frac{1}{2} H_{-1,0,0} - H_{-1,2} \Big] + (1-x) \left[9H_1 \zeta_2 \right. \\
& + 12H_{0,0,0,0} - \frac{293}{108} + \frac{61}{6} H_0 \zeta_2 - \frac{7}{3} H_{1,0} - \frac{857}{36} H_1 - 9H_0 \zeta_3 + 16H_{-2,-1,0} - 4H_{-2,0,0} + 8H_{-2} \zeta_2 \\
& - \frac{13}{2} H_{1,0,0} + \frac{3}{4} H_{1,1} - H_{1,1,0} - H_{1,1,1} \Big] + (1+x) \left[\frac{1}{6} H_{2,0} - \frac{95}{3} H_{-1,0} - \frac{149}{36} H_2 + \frac{3451}{108} H_0 \right. \\
& - 7H_{-2,0} + \frac{302}{9} H_{0,0} + \frac{19}{6} H_3 - \frac{991}{36} \zeta_2 - \frac{163}{6} \zeta_3 - \frac{35}{3} H_{0,0,0} + \frac{17}{6} H_{2,1} - \frac{43}{10} \zeta_2^2 + 13H_{-1} \zeta_2 \\
& + 18H_{-1,-1,0} - H_{3,1} - 6H_4 - 4H_{-1,2} + 6H_{0,0} \zeta_2 + 8H_2 \zeta_2 - 7H_{2,0,0} - 2H_{2,1,0} - 2H_{2,1,1} - 4H_{3,0} \\
& - 9H_{-1,0,0} \Big] - \frac{241}{288} \delta(1-x) + 16C_A n_f^2 \left(\frac{19}{54} H_0 - \frac{1}{24} x H_0 - \frac{1}{27} p_{gg}(x) + \frac{13}{54} \left(\frac{1}{x} - x^2 \right) \left[\frac{5}{3} - H_1 \right] \right. \\
& + (1-x) \left[\frac{11}{72} H_1 - \frac{71}{216} \right] + \frac{2}{9} (1+x) \left[\zeta_2 + \frac{13}{12} x H_0 - \frac{1}{2} H_{0,0} - H_2 \right] + \frac{29}{288} \delta(1-x) \Big) \\
& + 16C_A^2 n_f \left(x^2 \left[\zeta_3 + \frac{11}{9} \zeta_2 + \frac{11}{9} H_{0,0} - \frac{2}{3} H_3 + \frac{2}{3} H_0 \zeta_2 + \frac{1639}{108} H_0 - 2H_{-2,0} \right] + \frac{1}{3} p_{gg}(x) \left[\frac{10}{3} \zeta_2 \right. \right. \\
& - \frac{209}{36} - 8\zeta_3 - 2H_{-2,0} - \frac{1}{2} H_0 - \frac{10}{3} H_{0,0} - \frac{20}{3} H_{1,0} - H_{1,0,0} - \frac{20}{3} H_2 - H_3 \Big] + \frac{10}{9} p_{gg}(-x) \left[\zeta_2 \right. \\
& + 2H_{-1,0} + \frac{3}{10} H_0 \zeta_2 - H_{0,0} \Big] + \frac{1}{3} \left(\frac{1}{x} - x^2 \right) \left[H_3 - H_0 \zeta_2 - \frac{13}{3} H_2 + \frac{5443}{108} - 3H_1 \zeta_2 + \frac{205}{36} H_1 \right. \\
& - \frac{13}{3} H_{1,0} + H_{1,0,0} \Big] + \left(\frac{1}{x} + x^2 \right) \left[\frac{151}{54} H_0 - \frac{8}{3} \zeta_2 + \frac{1}{3} H_{-1} \zeta_2 - \zeta_3 + 2H_{-1,-1,0} - \frac{2}{3} H_{-1,0,0} \right. \\
& - \frac{37}{9} H_{-1,0} + \frac{2}{3} H_{-1,2} \Big] + (1-x) \left[\frac{5}{6} H_{-2,0} + H_{-3,0} + 2H_{0,0,0} - \frac{269}{36} \zeta_2 - \frac{4097}{216} - 3H_{-2} \zeta_2 \right. \\
& - 6H_{-2,-1,0} + 3H_{-2,0,0} - \frac{7}{2} H_1 \zeta_2 + \frac{677}{72} H_1 + H_{1,0} + \frac{7}{4} H_{1,0,0} \Big] + (1+x) \left[\frac{193}{36} H_2 - \frac{11}{2} H_{-1} \zeta_2 \right. \\
& + \frac{39}{20} \zeta_2^2 - \frac{7}{12} H_3 - \frac{53}{9} H_{0,0} + \frac{7}{12} H_0 \zeta_2 - \frac{5}{2} H_{0,0} \zeta_2 + 5\zeta_3 - 7H_{-1,-1,0} + \frac{77}{6} H_{-1,0} + \frac{9}{2} H_{-1,0,0} \\
& + 2H_{-1,2} - 3H_2 \zeta_2 - \frac{2}{3} H_{2,0} + \frac{3}{2} H_{2,0,0} + \frac{3}{2} H_4 \Big] + \frac{1}{9} \zeta_2 + 7H_{-2,0} + 2H_2 + \frac{458}{27} H_0 + H_{0,0} \zeta_2 \\
& + \frac{3}{2} \zeta_2^2 + 4H_{-3,0} - x \left[\frac{131}{12} H_{0,0} - \frac{8}{3} H_0 \zeta_2 + \frac{7}{2} H_3 - H_{0,0,0,0} + \frac{7}{6} H_{0,0,0} + \frac{1943}{216} H_0 + 6H_0 \zeta_3 \right] \\
& - 8(1-x) \left[\frac{233}{288} + \frac{1}{6} \zeta_2 + \frac{1}{12} \zeta_2^2 + \frac{5}{3} \zeta_3 \right] + 16C_A^3 \left(x^2 \left[33H_{-2,0} + 33H_0 \zeta_2 - \frac{1249}{18} H_0,0 \right. \right. \\
& - 44H_{0,0,0} - \frac{110}{3} H_3 - \frac{44}{3} H_{2,0} + \frac{85}{6} \zeta_2 + \frac{6409}{108} H_0 \Big] + p_{gg}(x) \left[\frac{245}{24} - \frac{67}{9} \zeta_2 - \frac{3}{10} \zeta_2^2 + \frac{11}{3} \zeta_3 \right]
\end{aligned}$$

$$\begin{aligned}
& - 4H_{-3,0} + 6H_{-2} \zeta_2 + 4H_{-2,-1,0} + \frac{11}{3} H_{-2,0} - 4H_{-2,0,0} - 4H_{-2,2} + \frac{1}{6} H_0 - 7H_0 \zeta_3 + \frac{67}{9} H_{0,0} \\
& - 8H_{0,0} \zeta_2 + 4H_{0,0,0,0} - 6H_1 \zeta_3 - 4H_{1,-2,0} + 10H_{2,0,0} - 6H_{1,0} \zeta_2 + 8H_{1,0,0,0} + 8H_{1,1,0,0} + 8H_4 \\
& + \frac{134}{9} H_{1,0} + \frac{11}{6} H_{1,0,0} + 8H_{1,2,0} + 8H_{1,3} + \frac{134}{9} H_2 - 4H_2 \zeta_2 + 8H_{3,1} + 8H_{2,2} + \frac{11}{6} H_3 + 10H_{3,0} \\
& + 8H_{2,1,0} \Big] + p_{gg}(-x) \left[\frac{11}{2} \zeta_2^2 - \frac{11}{6} H_0 \zeta_2 - 4H_{-3,0} + 16H_{-2} \zeta_2 - 12H_{-2,2} - \frac{134}{9} H_{-1,0} + 2H_2 \zeta_2 \right. \\
& + 8H_{-2,-1,0} + 12H_{-1} \zeta_3 - 18H_{-2,0,0} + 8H_{-1,-2,0} - 16H_{-1,-1} \zeta_2 + 24H_{-1,-1,0,0} + 16H_{-1,-1,2} \\
& + 18H_{-1,0} \zeta_2 - 16H_{-1,0,0,0} - 4H_{-1,2,0} - 16H_{-1,3} - 5H_0 \zeta_3 - 8H_{0,0} \zeta_2 + 4H_{0,0,0,0} + 2H_{3,0} \\
& - \frac{67}{9} \zeta_2 + \frac{67}{9} H_{0,0} + 8H_4 \Big] + \left(\frac{1}{x} - x^2 \right) \left[\frac{16619}{162} + \frac{22}{3} H_{2,0} - \frac{55}{2} \zeta_3 - \frac{11}{2} H_0 \zeta_2 - \frac{67}{9} H_2 - \frac{67}{9} H_{1,0} \right. \\
& - \frac{413}{108} H_1 - \frac{11}{2} H_1 \zeta_2 + \frac{33}{2} H_{1,0,0} \Big] + 11 \left(\frac{1}{x} + x^2 \right) \left[\frac{71}{54} H_0 - \frac{1}{6} H_3 - \frac{389}{198} \zeta_2 - \frac{2}{3} H_{-2,0} - \frac{1}{2} H_{-1} \zeta_2 \right. \\
& + H_{-1,-1,0} - \frac{523}{198} H_{-1,0} + \frac{8}{3} H_{-1,0,0} + H_{-1,2} \Big] + (1-x) \left[\frac{31}{36} H_1 + \frac{27}{2} H_{1,0} - \frac{25}{2} H_{1,0,0} - 4H_{-3,0} \right. \\
& - \frac{263}{12} H_{0,0} - \frac{29}{3} H_{0,0,0} - \frac{19}{3} H_{-2,0} - \frac{11317}{108} - 4H_{-2} \zeta_2 - 8H_{-2,-1,0} - 12H_{-2,0,0} - \frac{3}{2} H_1 \zeta_2 \Big] \\
& + (1+x) \left[\frac{27}{2} H_0 \zeta_2 - \frac{43}{6} H_3 + \frac{29}{3} H_{2,0} + \frac{4651}{216} H_0 - \frac{329}{18} \zeta_2 + \frac{11}{2} (1+x) \zeta_3 - \frac{43}{5} \zeta_2^2 - \frac{215}{6} H_{-1,0} \right. \\
& - 22H_{0,0} \zeta_2 - 8H_0 \zeta_3 - 3H_{-1,-1,0} + 38H_{-1,0,0} + 25H_{-1,2,0} + 10H_{2,0,0} - 4H_2 \zeta_2 + 16H_{3,0} + 26H_4 \\
& - \frac{158}{9} H_2 - \frac{53}{2} H_{-1} \zeta_2 \Big] - 29H_{0,0} - \frac{40}{3} H_{0,0,0} + 27H_{-2,0} + \frac{41}{3} H_0 \zeta_2 - 20H_3 - 24H_{2,0} + \frac{53}{6} \zeta_2 \\
& + \frac{601}{12} H_0 + 24\zeta_3 + 2\zeta_2^2 + 27H_2 - 4H_0 \zeta_2 - 16H_0 \zeta_3 - 16H_{-3,0} + 28xH_{0,0,0,0} + \delta(1-x) \left[\frac{79}{32} \right. \\
& - \zeta_2 \zeta_3 + \frac{1}{6} \zeta_2 + \frac{11}{24} \zeta_2^2 + \frac{67}{6} \zeta_3 - 5\zeta_5 \Big] + 16C_F n_f^2 \left(\frac{2}{9} x^2 \left[\frac{11}{6} H_0 + H_2 - \zeta_2 + 2H_{0,0} - 9 \right] + \frac{1}{3} H_1 \right. \\
& - \frac{1}{3} \zeta_2 - \frac{10}{3} H_0 - \frac{1}{3} H_{0,0} + 2 + \frac{2}{9} \left(\frac{1}{x} - x^2 \right) \left[\frac{8}{3} H_1 - 2H_{1,0} - H_{1,1} - \frac{77}{18} \right] - (1-x) \left[\frac{1}{3} H_{1,0} + \frac{1}{6} H_{1,1} \right. \\
& + \frac{4}{9} + \frac{13}{6} H_1 + xH_1 \Big] + \frac{1}{3} (1+x) \left[\frac{68}{9} H_0 - \frac{4}{3} H_2 + \frac{4}{3} \zeta_2 + \frac{29}{6} H_{0,0} - \zeta_3 + 2H_0 \zeta_2 - H_{0,0,0} - 2H_3 \right. \\
& - H_{2,1} - 2H_{2,0} \Big] + \frac{11}{144} \delta(1-x) \Big) + 16C_F^2 n_f \left(\frac{4}{3} x^2 \left[\frac{163}{16} + \frac{27}{8} H_0 + \frac{7}{2} H_{0,0} - H_{2,0} - \zeta_2 + \frac{9}{4} H_{1,0} \right. \right. \\
& - H_{2,1} + \frac{1}{2} H_{0,0,0} + \frac{85}{16} H_1 + H_2 - 2H_{-2,0} - \frac{3}{2} \zeta_3 \Big] + \frac{4}{3} \left(\frac{1}{x} - x^2 \right) \left[\frac{31}{16} H_1 - \frac{11}{16} - \frac{5}{4} H_{1,0} + \frac{1}{2} H_{1,0,0} \right. \\
& - H_1 \zeta_2 - H_{1,1} + H_{1,1,0} + H_{1,1,1} + \zeta_3 \Big] + \frac{4}{3} \left(\frac{1}{x} + x^2 \right) \left[H_{-1} \zeta_2 + 2H_{-1,-1,0} - H_{-1,0,0} \right] + \frac{215}{12} H_{0,0} \\
& + \frac{20}{3} H_0 - \frac{131}{6} + 3H_{2,0} + \frac{205}{12} x \zeta_2 - 3H_{1,0} + H_{2,1} - \frac{85}{12} H_1 + \frac{11}{4} H_2 + 8H_{-2,0} + 2\zeta_2^2 - H_0 \zeta_2 \\
& + H_3 + 6H_0 \zeta_3 + 8H_{-3,0} - 4xH_{0,0,0} + (1-x) \left[\frac{107}{12} H_1 - \frac{5}{6} H_{1,0} - 4\zeta_2 + H_0 \zeta_3 - 8H_{-2,-1,0} \right. \\
& - 4H_{-2} \zeta_2 + 4H_{-2,0,0} - 4H_1 \zeta_2 + \frac{7}{2} H_{1,0,0} - \frac{7}{12} H_{1,1} + H_{1,1,0} + H_{1,1,1} \Big] + (1+x) \left[\frac{5}{4} H_2 + \frac{33}{8} \right. \\
& - \frac{99}{4} H_{0,0} - 8H_{2,0} - \frac{541}{24} H_0 - 4H_{2,1} - \frac{3}{2} H_{0,0,0} - 2x \zeta_3 + \frac{9}{2} \zeta_2^2 + 5H_0 \zeta_2 - 5H_3 - 4H_{-1} \zeta_2 \\
& - 8H_{-1,-1,0} + \frac{67}{3} H_{-1,0} + 4H_{-1,0,0} + 2H_0 \zeta_2 - 2H_{0,0,0,0} - 4H_2 \zeta_2 + 3H_{2,0,0} + 2H_{2,1,0} \\
& + 2H_{2,1,1} + H_{3,1} - 2H_4 \Big] + \frac{1}{16} \delta(1-x)
\end{aligned}$$



Large- x behaviour

Structure at three loops

$$P_{aa,x \rightarrow 1}^{(2)}(x) = \frac{A_3^a}{(1-x)_+} + B_3^a \delta(1-x) + C_3^a \ln(1-x) + \mathcal{O}(1)$$

Korchemsky (89)

Leading terms \rightarrow soft-gluon resummation

$$\begin{aligned} A_3^q &= +16 C_F C_A^2 \left(\frac{245}{24} - \frac{67}{9} \zeta_2 + \frac{11}{6} \zeta_3 + \frac{11}{5} \zeta_2^2 \right) \\ &\quad + 16 C_F C_A n_f \left(-\frac{209}{108} + \frac{10}{9} \zeta_2 - \frac{7}{3} \zeta_3 \right) \\ &\quad + 16 C_F^2 n_f \left(-\frac{55}{24} + 2 \zeta_3 \right) + 16 C_F n_f^2 \left(-\frac{1}{27} \right) \end{aligned}$$

$C_A A_3^q = C_F A_3^g$. n_f part of A_3^q also by C. Berger (02)

n_f^2 part Gracey (94)

Subleading logarithms: surprising relation

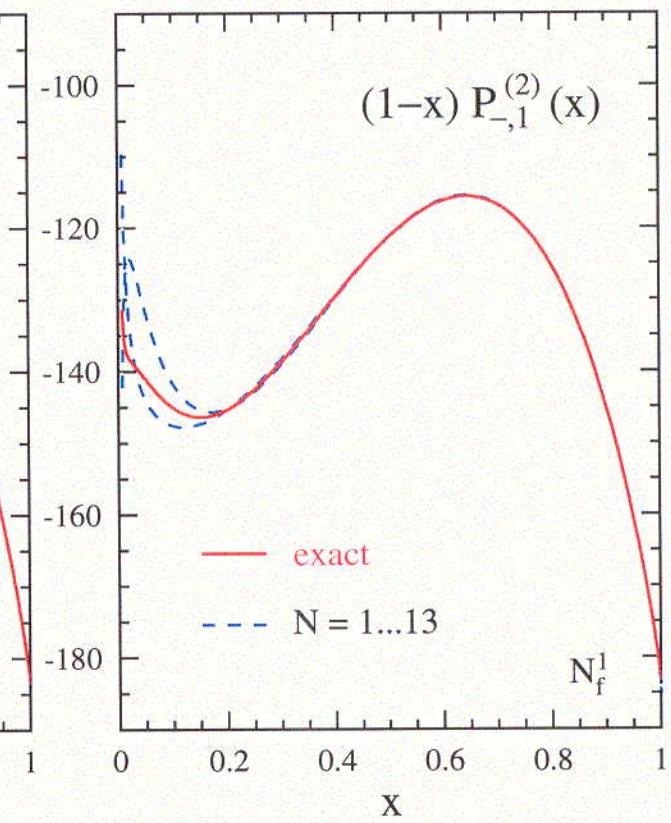
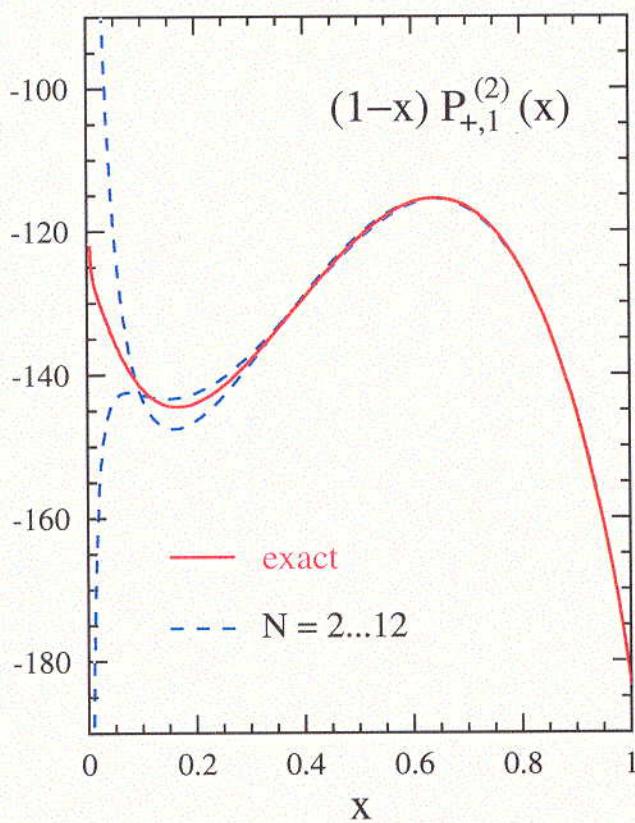
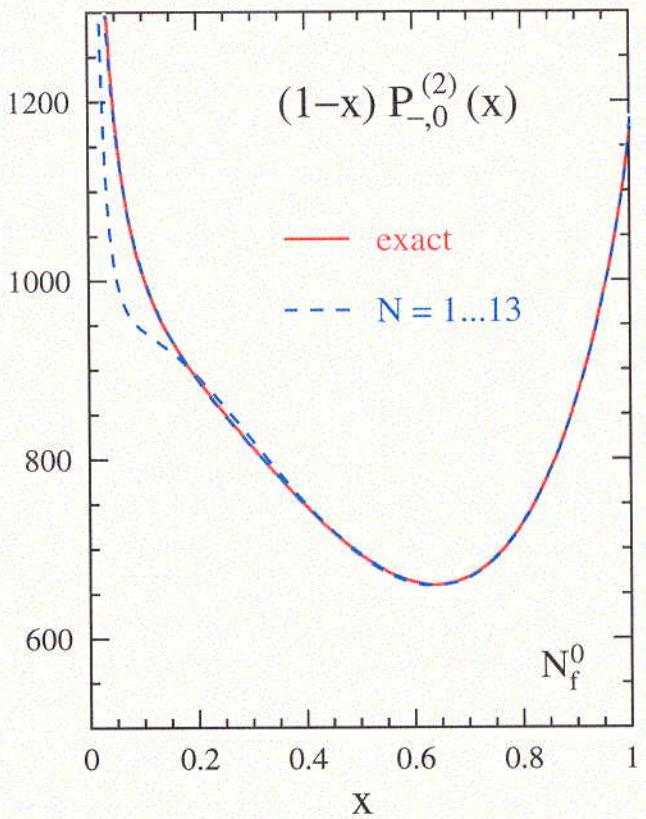
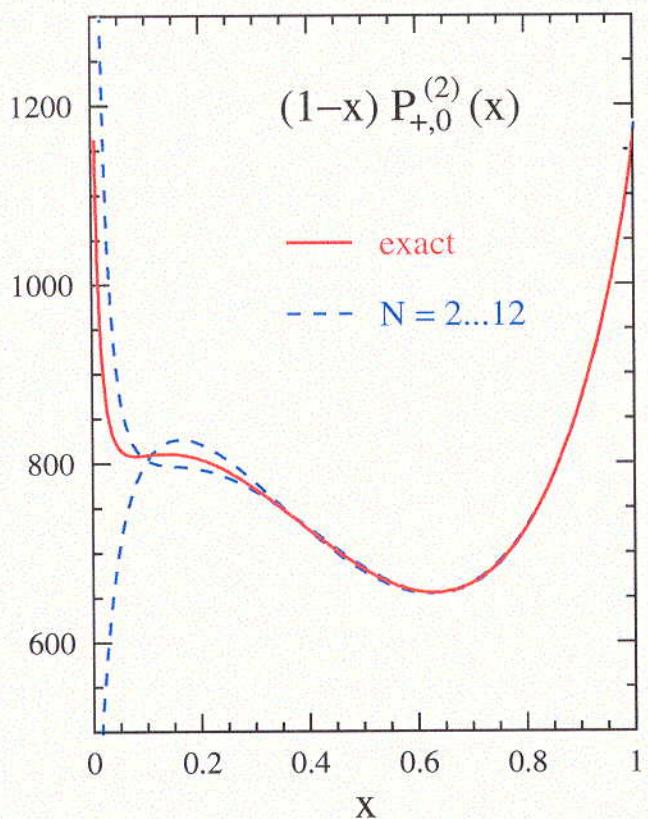
$$\begin{aligned} C_1^q &= 0, & C_2^q &= 4 C_F A_1^q, & C_3^q &= 8 C_F A_2^q \\ C_1^g &= 0, & C_2^g &= 4 C_A A_1^g, & C_3^g &= 8 C_A A_2^g \end{aligned}$$

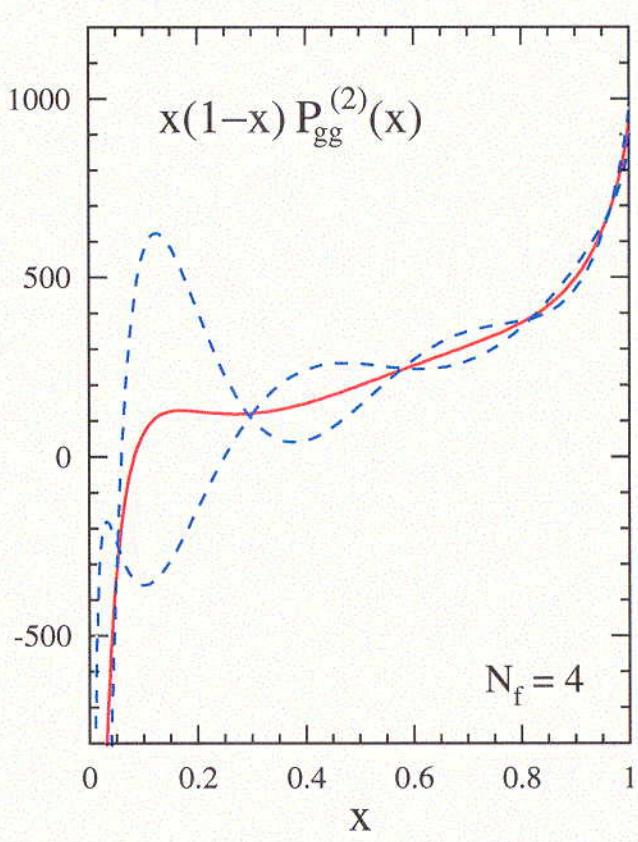
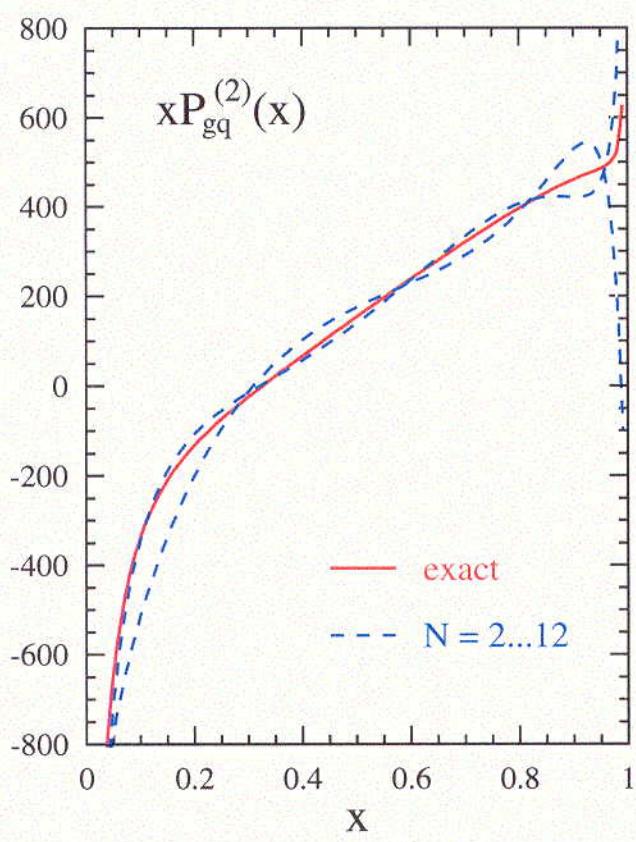
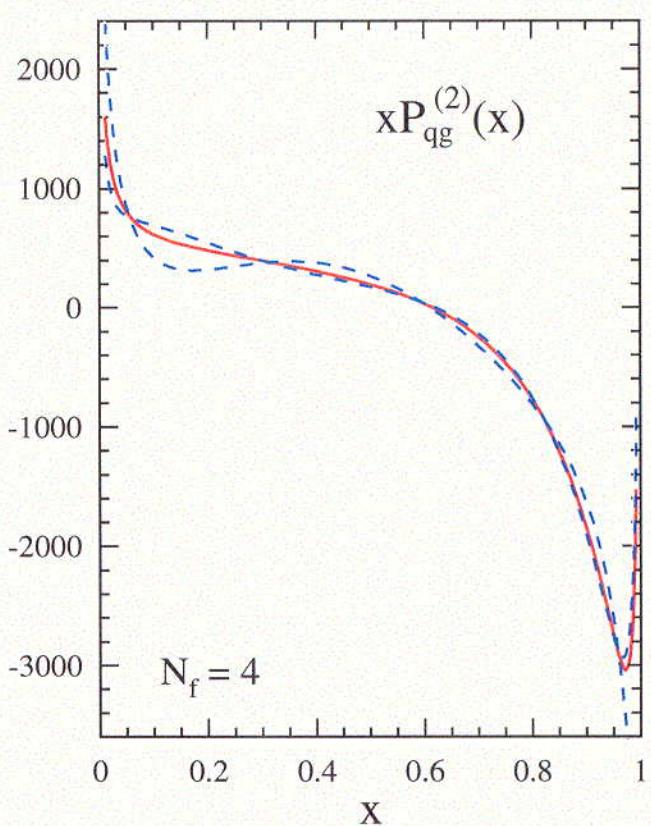
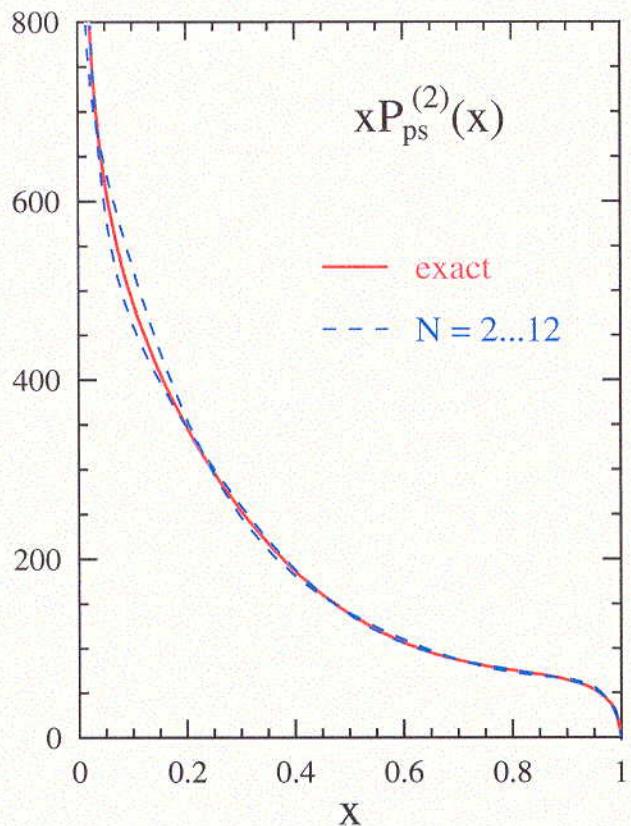
Suggests yet unexplored structure

Off-diagonal contributions

$$P_{ab,x \rightarrow 1}^{(2)}(x) = \sum_{i=0}^3 D_i^{ab} \ln^{4-i}(1-x) + \mathcal{O}(1)$$

$D_{0,1,2}^{ab} = 0$ for $C_A \equiv N_c = C_F = n_f$: SUSY relation





Small- x : non-singlet case

Structure at three loops

$$P_{x \rightarrow 0}^{(2)i}(x) = D_0^i \ln^4 x + \dots + D_3^i \ln x + \mathcal{O}(1)$$

Generally terms up to $\ln^{2k} x$ at order α_s^{k+1}

Coefficients for ‘plus’ and ‘minus’ cases

$$D_0^+ \cong 1.58025$$

$$D_1^+ \cong 29.6296 - 2.37037 n_f$$

$$D_2^+ \cong 295.042 - 32.1975 n_f + 0.592592 n_f^2$$

$$D_3^+ \cong 1261.11 - 152.597 n_f + 4.345679 n_f^2$$

$$D_0^- \cong 1.43210$$

$$D_1^- \cong 35.5556 - 3.16049 n_f$$

$$D_2^- \cong 399.205 - 39.7037 n_f + 0.592592 n_f^2$$

$$D_3^- \cong 1465.93 - 172.693 n_f + 4.345679 n_f^2$$

D_0^\pm : Kirschner, Lipatov (83), Blümlein, A.V. (95)

Large logarithms often have small coefficients

New $d_{abc}d_{abc}$ contribution to ‘valence’ case

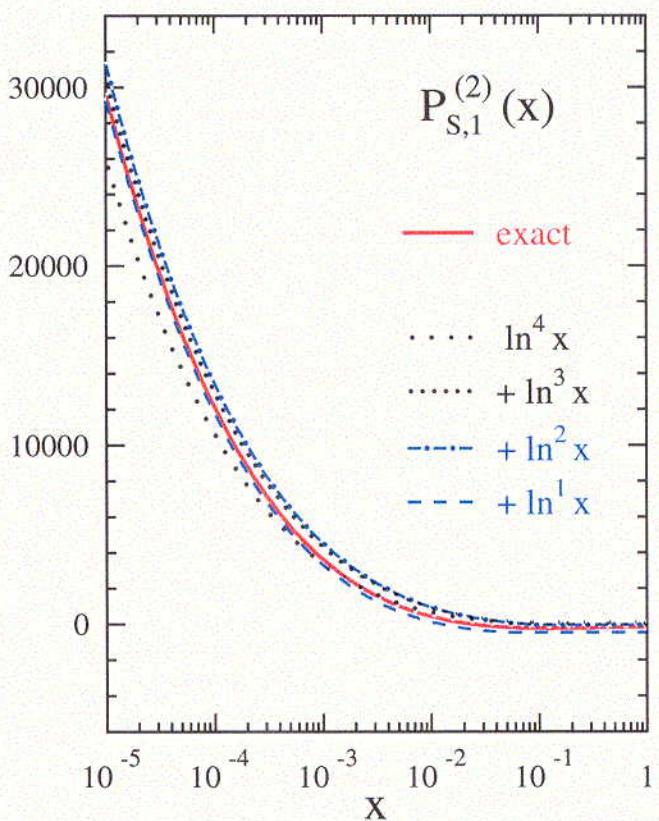
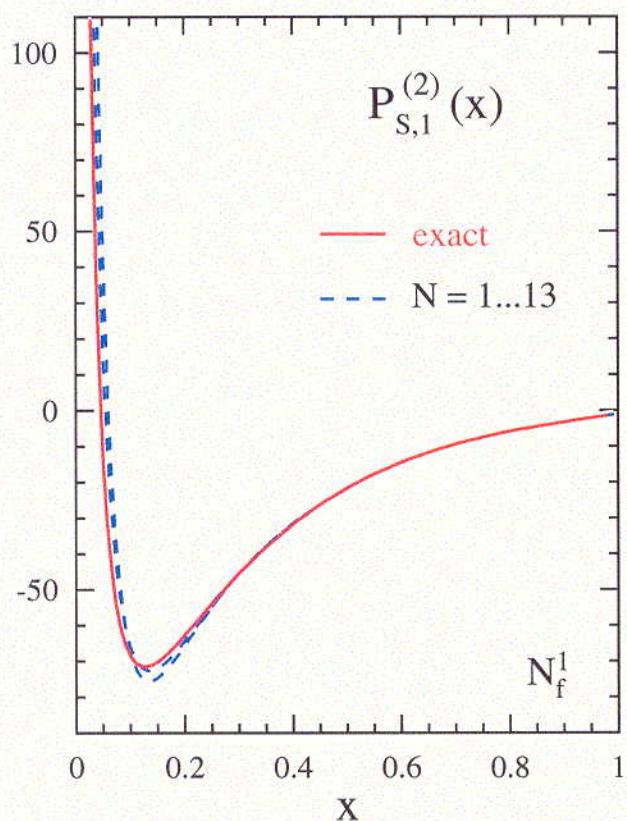
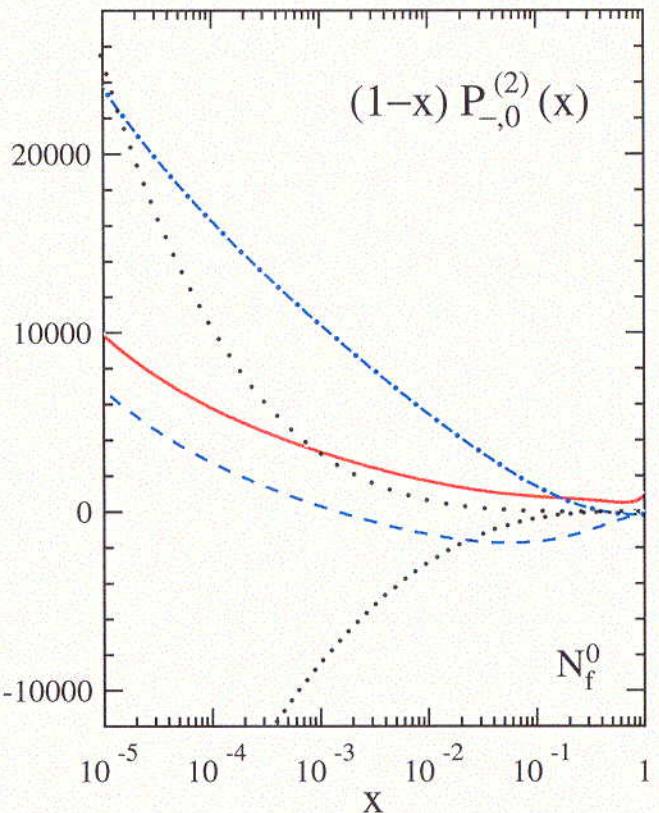
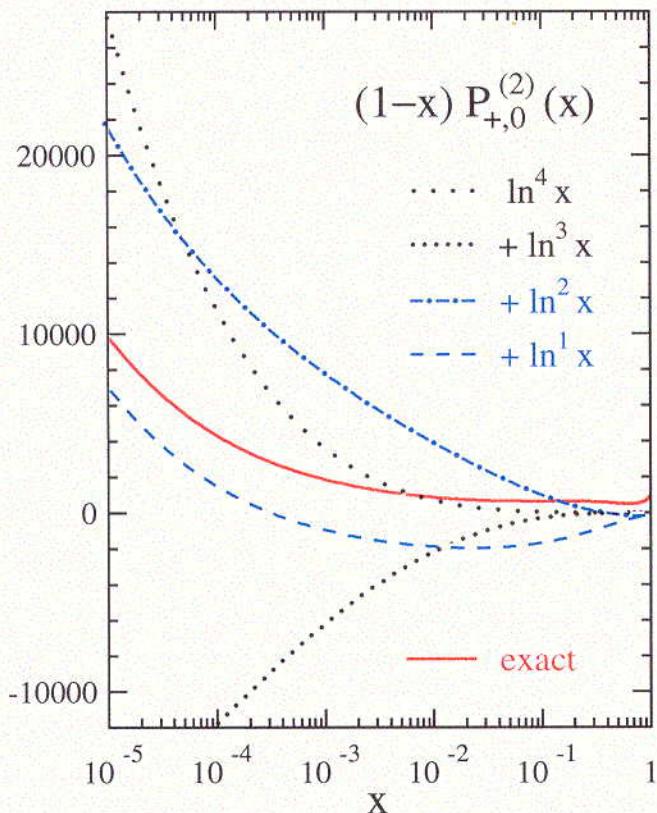
$$D_0^S \cong +1.48148 n_f$$

$$D_1^S \cong -2.96296 n_f$$

$$D_2^S \cong +6.89182 n_f$$

$$D_3^S \cong +178.030 n_f$$

Presence of $\ln^4 x$ term unpredicted



Small- x : flavour-singlet case

Structure at three loops

$$P_{ab,x \rightarrow 0}^{(2)}(x) = E_1^{ab} \frac{\ln x}{x} + E_2^{ab} \frac{1}{x} + \mathcal{O}(\ln^4 x)$$

Generally terms up to $x^{-1} \ln^k x$ at order α_s^{k+1}

Coefficients of $1/x$ contributions

$$E_1^{qq} \approx -132.741 n_f$$

$$E_2^{qq} \approx -505.999 n_f + 3.16049 n_f^2$$

$$E_1^{qg} \approx -298.667 n_f$$

$$E_2^{qg} \approx -1268.28 n_f + 4.57613 n_f^2$$

$$E_1^{gg} \approx +1189.27 + 71.0825 n_f$$

$$E_2^{gg} \approx +6163.11 - 46.4075 n_f - 2.37037 n_f^2$$

$$E_1^{gg} \approx +2675.85 + 157.269 n_f$$

$$E_2^{gg} \approx +14214.2 + 182.958 n_f - 2.79835 n_f^2$$

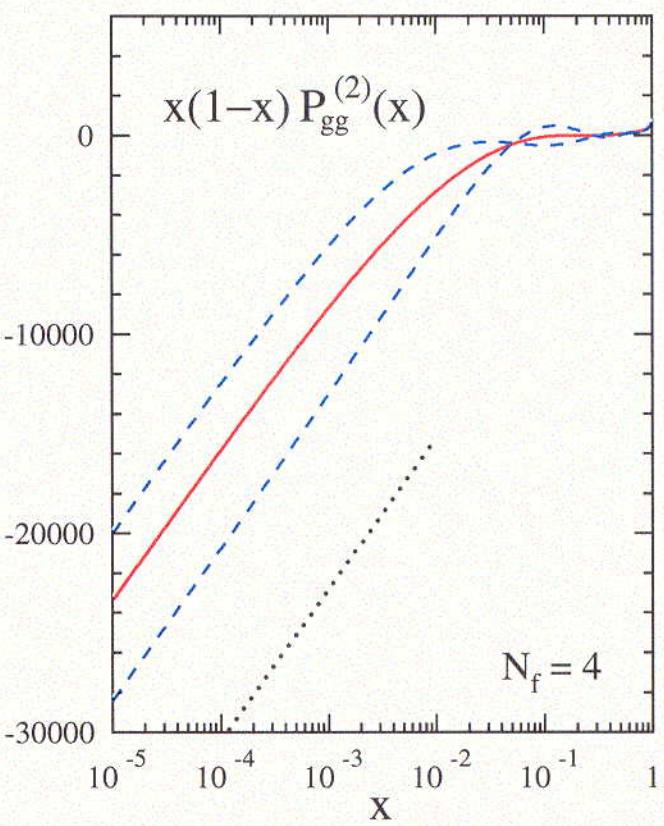
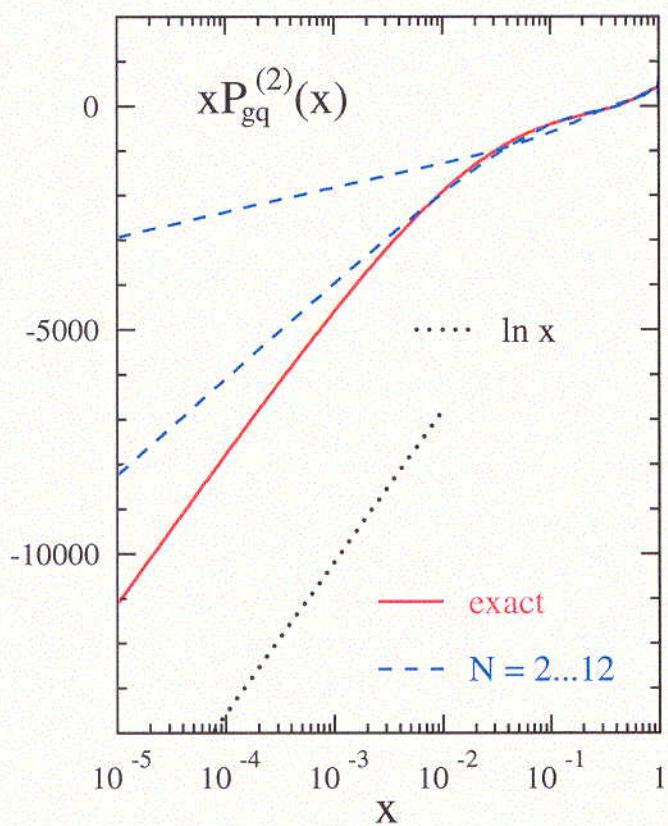
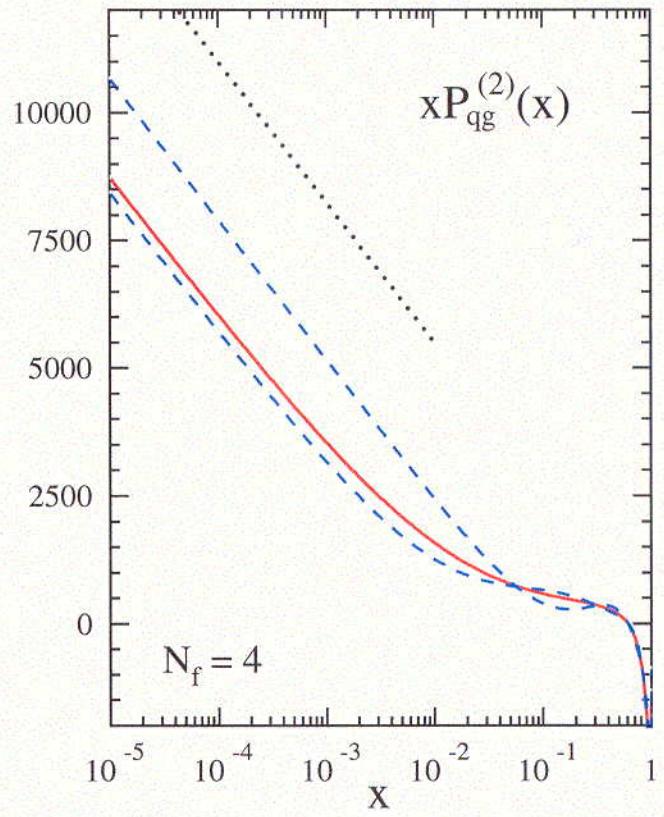
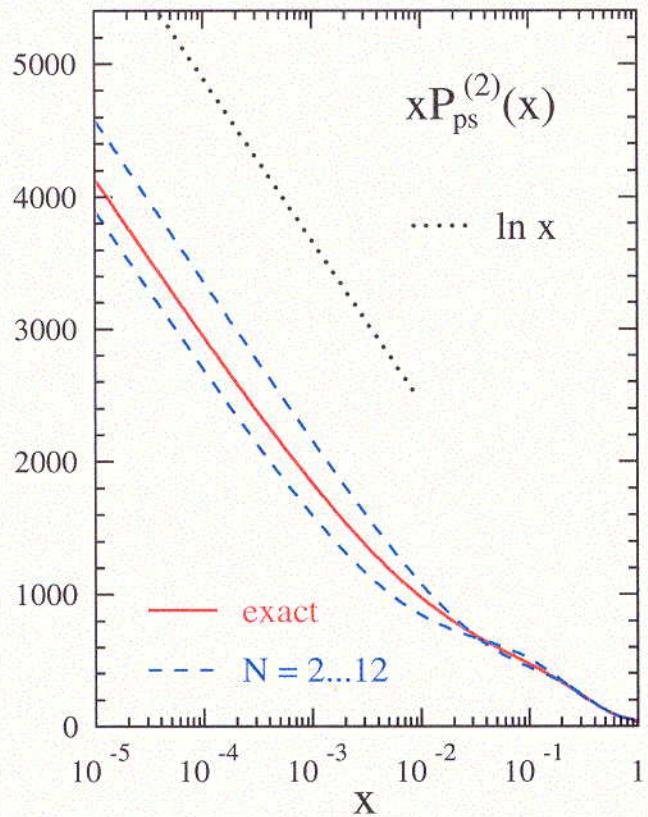
E_1^{qa} Catani, Hautmann (94), E_1^{gg} Fadin, Lipatov (98)

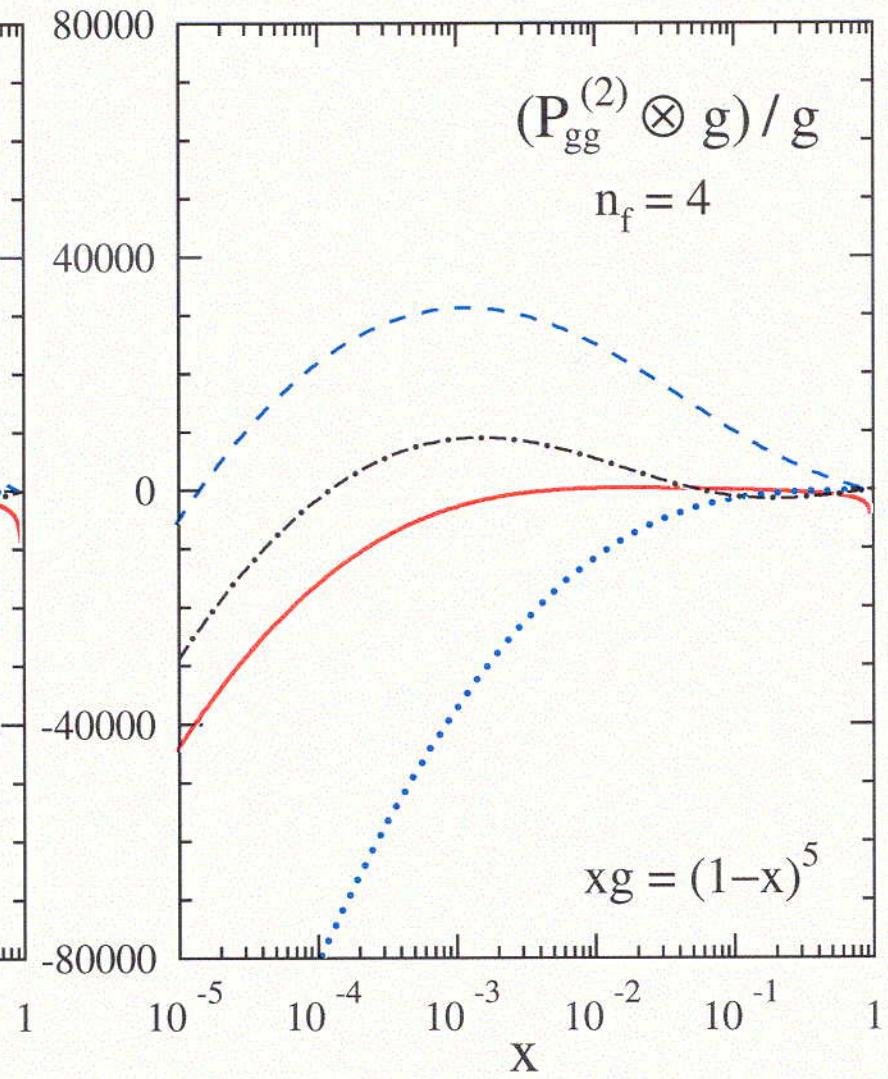
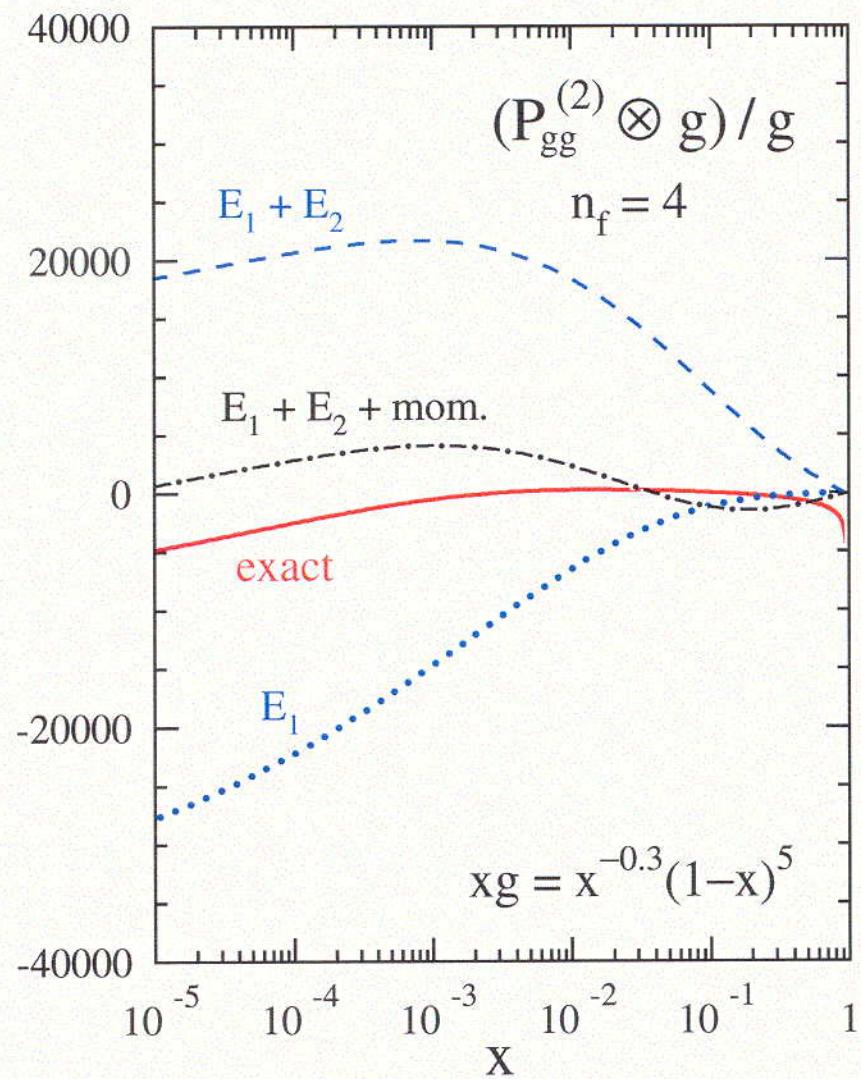
E_2^{ab}/E_1^{ab} for $n_f = 3 \dots 5$: 3.7 ... 4.7 Cf. $\ln 10^4 \simeq 9.2$

Relations between coefficients

$$E_1^{qg} = \frac{C_A}{C_F} E_1^{qq}$$

$$E_1^{gg} = \frac{C_A}{C_F} E_1^{qg} - \frac{8}{3} n_f$$





Parton evolution at NNLO

Illustrated by logarithmic derivatives $\dot{f} = d \ln f / d \ln \mu_f^2$

Order-independent model distributions

$$xq_{\text{ns}}(x, \mu_0^2) = x^{0.5} (1-x)^3$$

$$xq_s(x, \mu_0^2) = 0.6 x^{-0.3} (1-x)^{3.5} (1 + 5.0 x^{0.8})$$

$$xg(x, \mu_0^2) = 1.6 x^{-0.3} (1-x)^{4.5} (1 - 0.6 x^{0.3})$$

for

$$\alpha_s(\mu_0^2) = 0.2, \quad n_f = 4$$

Renormalization scale dependence

$$P_{ab}^{\text{NNLO}}(\mu_f, \mu_r) =$$

$$a_s(\mu_r^2) P_{ab}^{(0)}$$

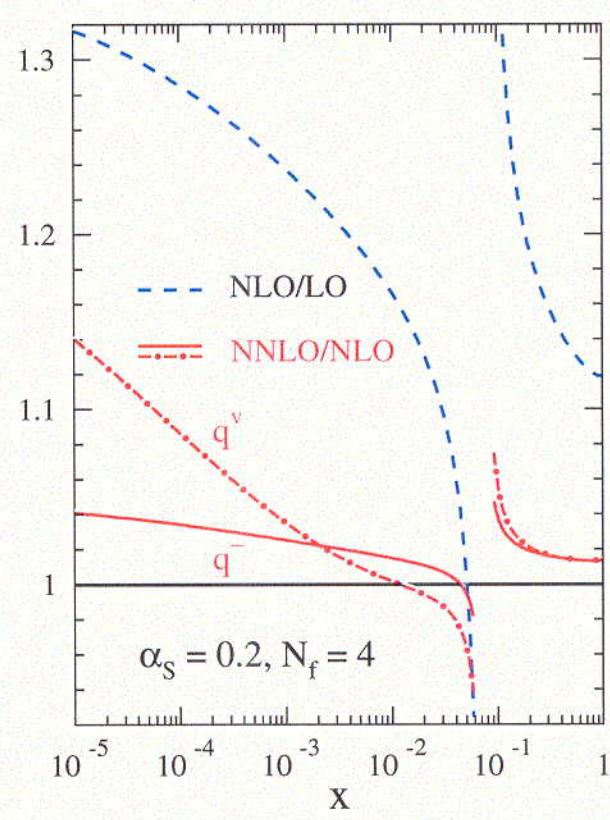
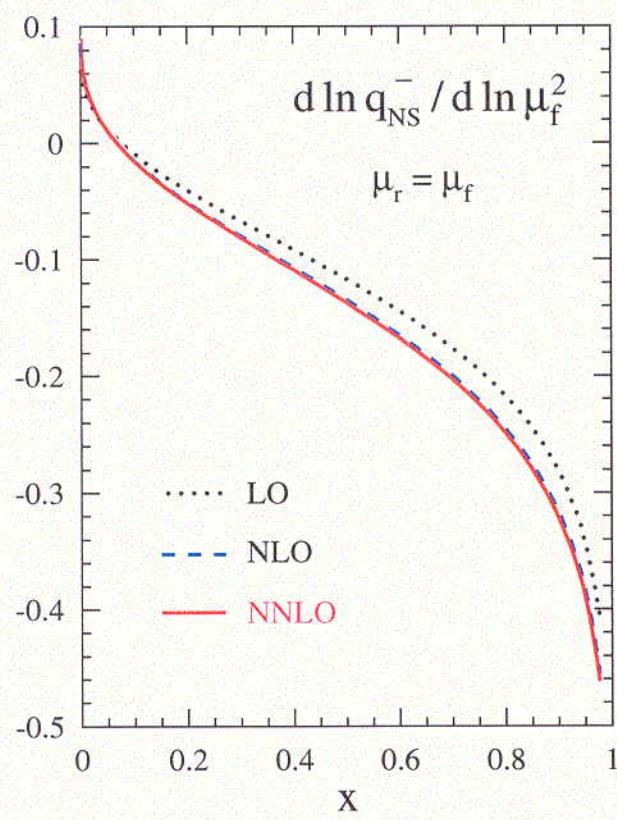
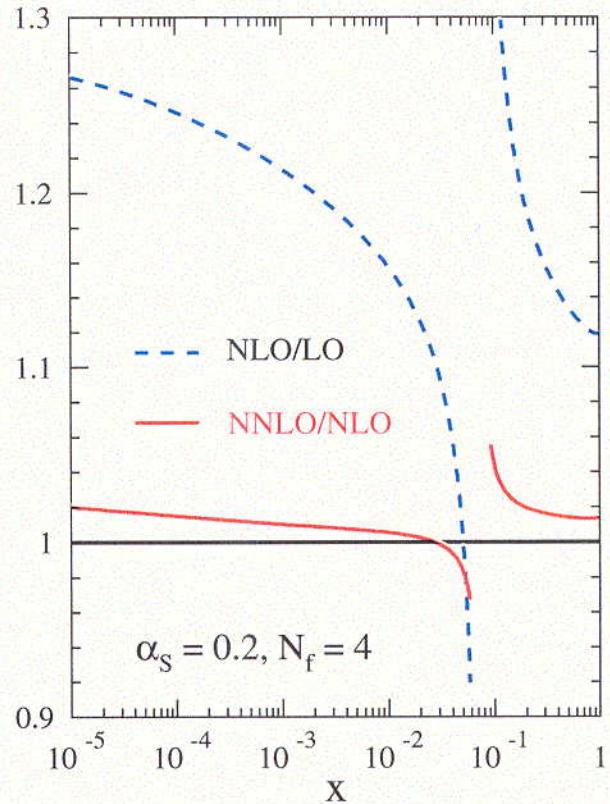
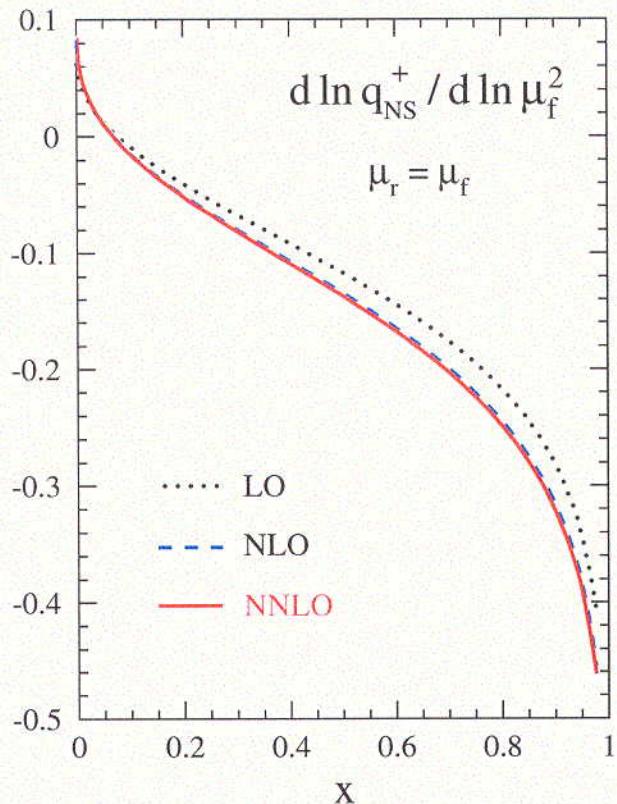
$$+ a_s^2(\mu_r^2) \left(P_{ab}^{(1)} - \beta_0 P_{ab}^{(0)} L \right)$$

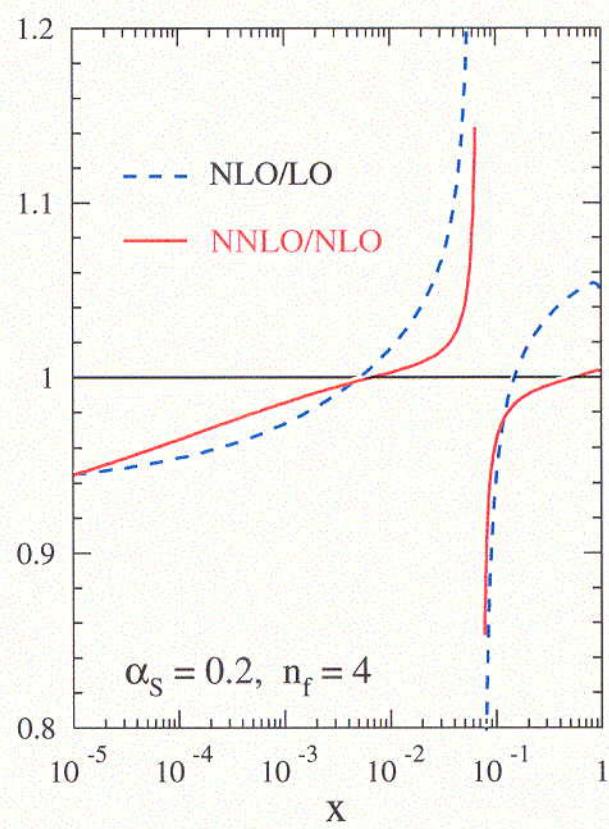
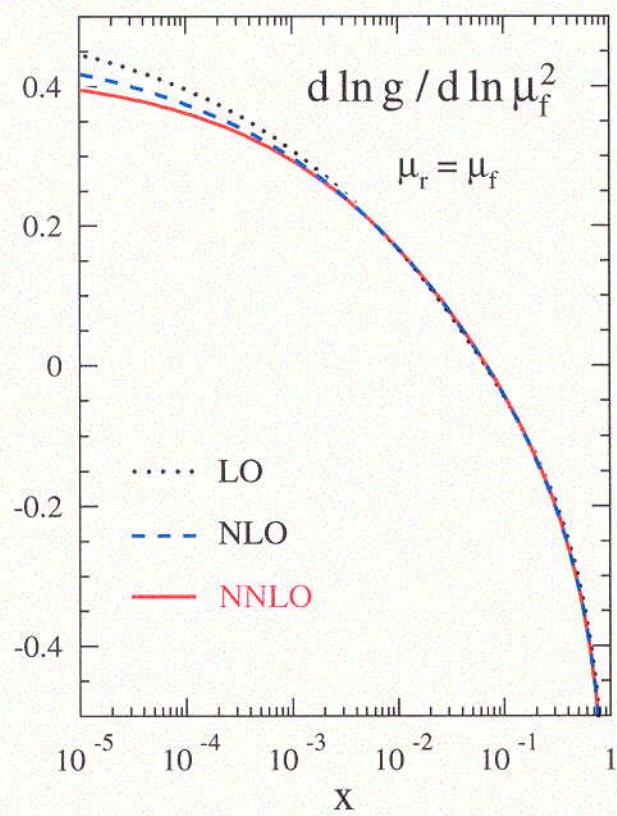
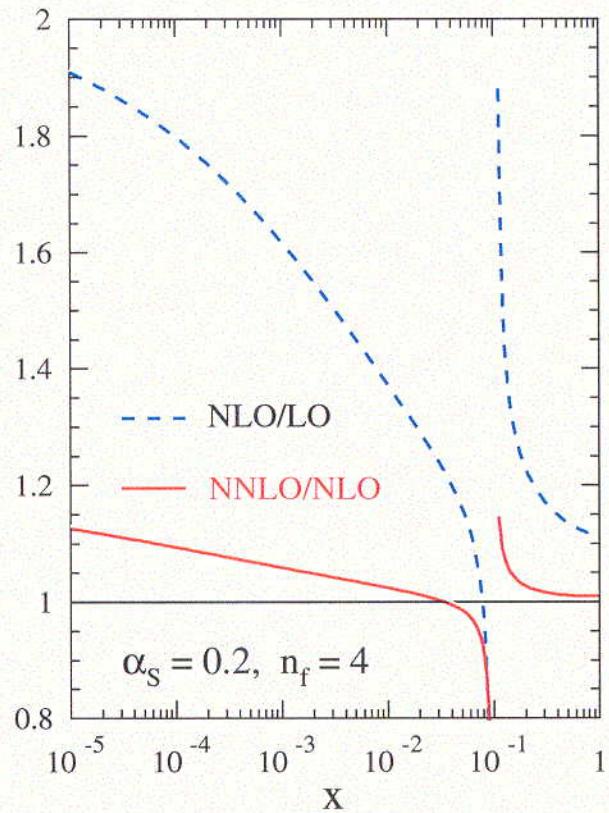
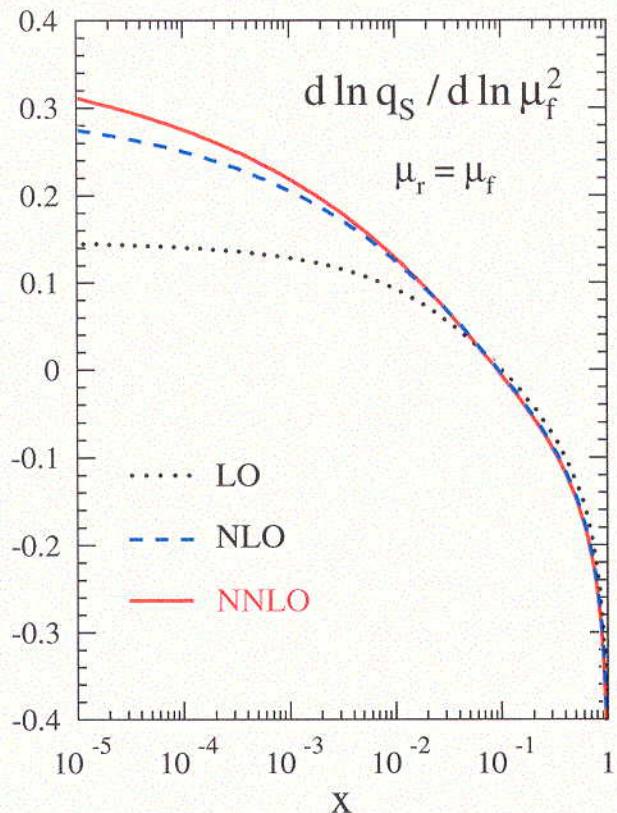
$$+ a_s^3(\mu_r^2) \left(P_{ab}^{(2)} - \left\{ \beta_1 P_{ab}^{(0)} + 2\beta_0 P_{ab}^{(1)} \right\} L + \beta_0^2 P_{ab}^{(0)} L^2 \right)$$

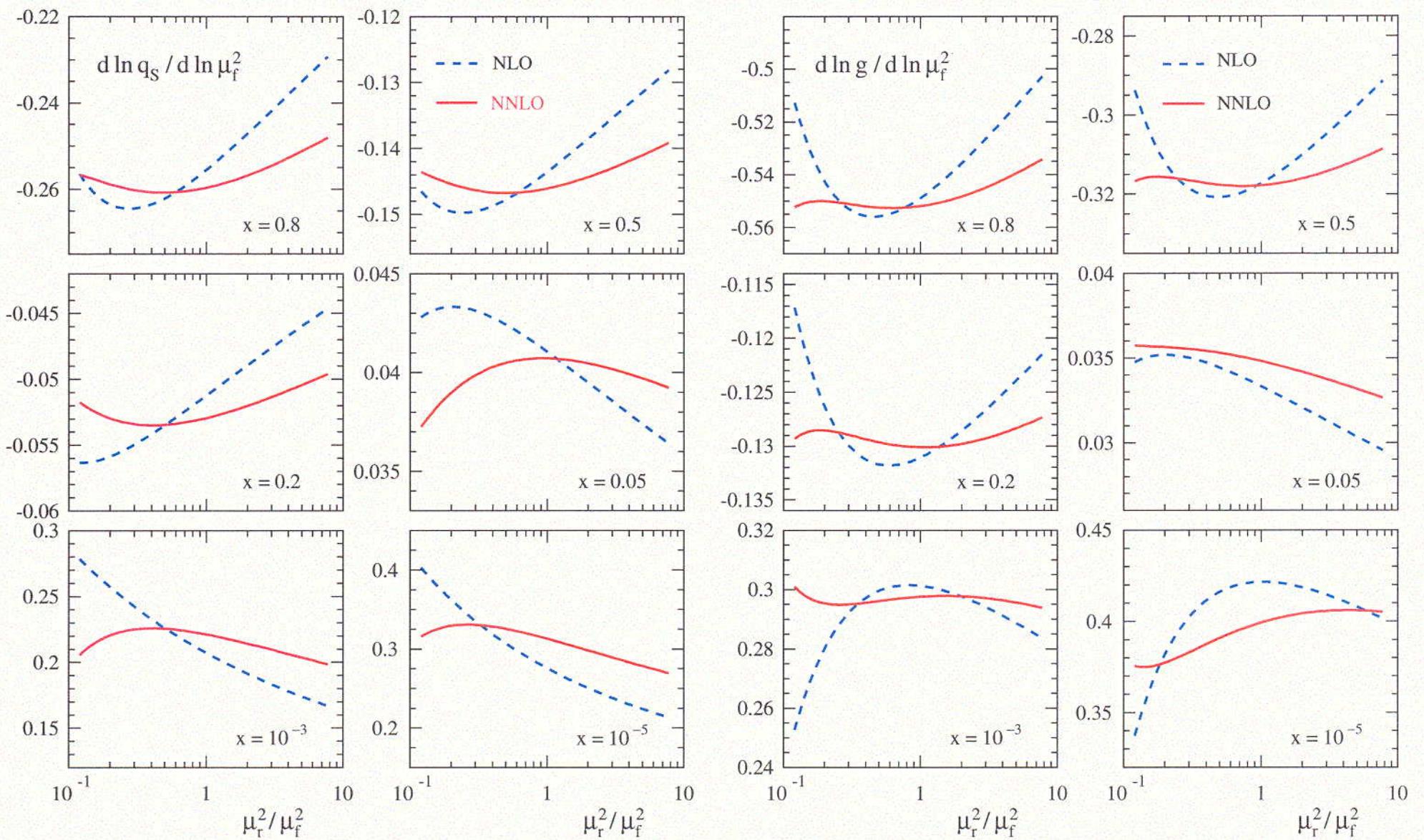
Abbreviations: $a_s = \alpha_s/(4\pi)$, $L = \ln(\mu_f^2/\mu_r^2)$

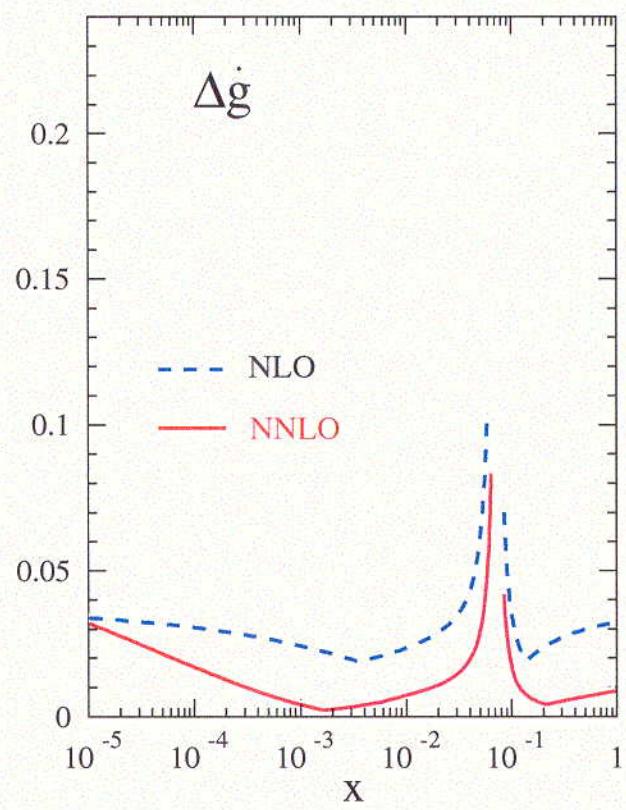
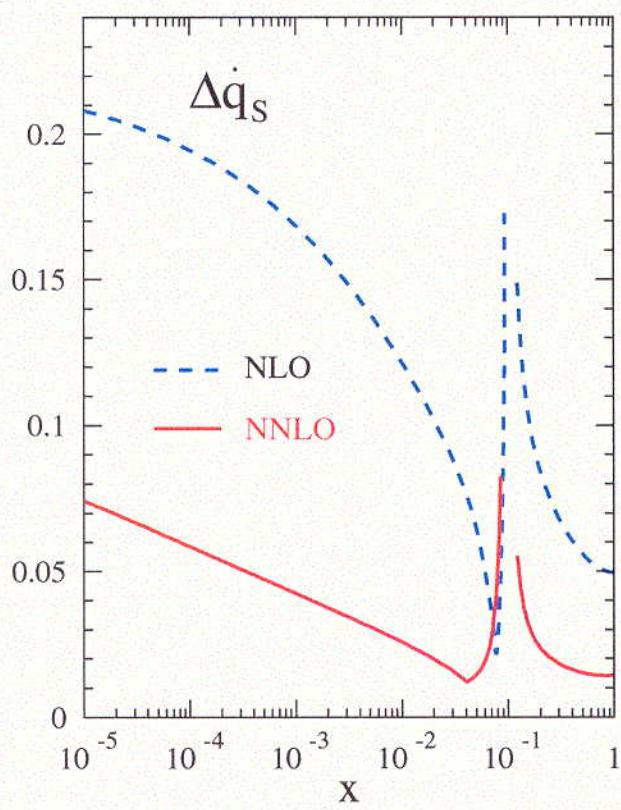
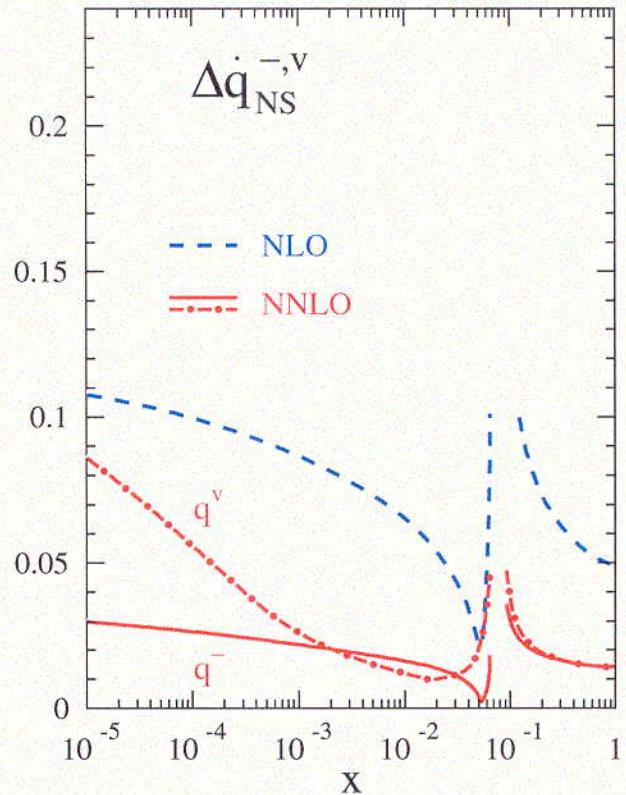
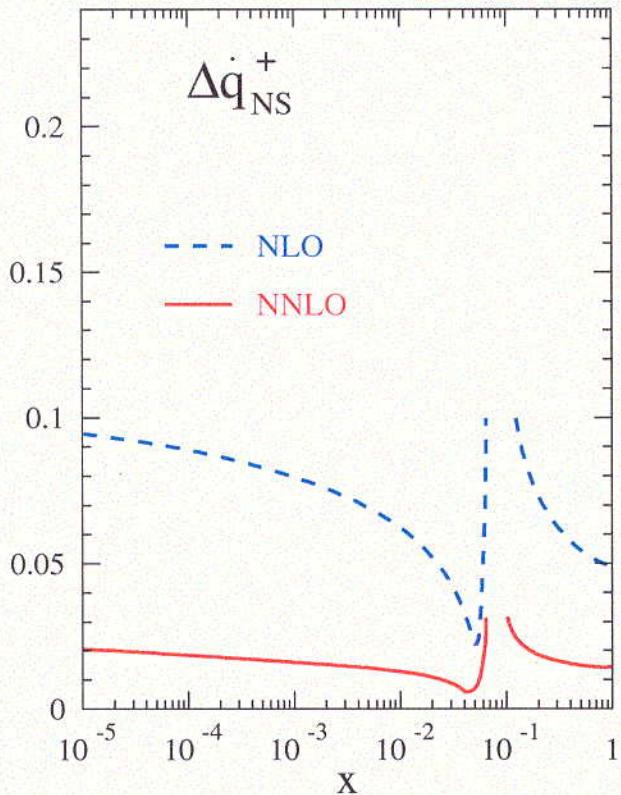
Conventional estimate for relative scale uncertainty

$$\Delta \dot{f} = \frac{\max \dot{f} - \min \dot{f}}{2 |\text{average } \dot{f}|}, \quad \mu_r^2 = \frac{1}{4} \mu_f^2 \dots 4 \mu_f^2$$









Summary

Unpolarized three-loop (NNLO) splitting functions calculated via all- N expressions in moment space

- Check of the Form manipulations at all stages
Fall back to small fixed $N \rightarrow$ Mincer program
Gorishny et al. (1989); Larin et al. (91)
- Barely possible with present resources
 ~ 10 person-years, several CPU years, Form 3.1
 10^4 diagrams, $\sim 10^5$ tabulated integrals (> 3 GB)
- To come: three-loop coefficient functions
Planned : polarized case, photon structure

Agreement with all known partial results

- Fixed low-integer moments
Larin et al. (94,97); Retey, Vermaseren (00)
- Large- n_f limits
Gracey (94); Bennett, Gracey (97)
- Large- x structure and limits
Korchemsky (89), C. Berger (02)
- Small- x limits
Catani, Hautmann (94)
Blümlein, A.V. (95)
Fadin, Lipatov (98)
- Approximations
van Neerven, A.V. (00)

Summary

NNLO effects on PDF evolution small at $x \gtrsim 10^{-3}$

Large x for $\alpha_S = 0.2$: corrections, μ_r variation: < 2%,
3-loop effects ~ 8 times smaller than 2-loop corr.

Evolution perturbative down to rather large α_S

Small- x : corrections increase with decreasing x

Three lessons at small momentum fractions

- New colour factors can give dominant higher-order terms not predicted by resummations
- The leading terms for $x \rightarrow 0$ do not dominate the splitting functions at exp. relevant x -values
- Convolutions $P \otimes f$ entering the evolution are sensitive even to terms power-suppressed in x

Small- x constraints need to be complemented by large- x information: at least several moments

Relations found between diagonal large- x terms

Coeff. of $\ln(1-x) \sim$ coeff. of $1/(1-x)$ at order $n-1$